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# Microkinetic modeling of the NO +  $H_2$  system on Pt/Al<sub>2</sub>O<sub>3</sub> catalyst using temporal analysis of products

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#### **ABSTRACT**

A systematic experimental and modeling study of the NO +  $H_2$  reaction system on Pt/Al<sub>2</sub>O<sub>3</sub> catalyst was carried out using temporal analysis of products (TAP). A microkinetic model was developed that explains the transient product distribution during various experimental TAP protocols. NO pulsing experiments show the inhibiting effect of oxygen poisoning and the kinetic limitations posed by N–O bond scission. The NO–H<sub>2</sub> pump–probe experiments demonstrate the effect of temperature, H<sub>2</sub>/NO ratio ( $\geq$ 1), and pulse delay time between consecutive NO and  $H_2$  pulses on  $N_2$ ,  $N_2O$ , and  $NH_3$  selectivity. The simulations reveal the competition between surface N–N recombination and N and/or NO reaction with H to form  $N_2$ and NH3, respectively. The developed mechanistic-based microkinetic model accounts for all of the experimental data including the aforementioned selectivity versus pulse delay trends.

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#### 1. Introduction

 $NO<sub>x</sub>$  storage and reduction (NSR) is an emerging technology for  $NO<sub>x</sub>$  emission abatement in lean burn gasoline and diesel engines. NSR consists of two stages of operations: storage and reduction. The first stage involves storage of  $NO<sub>x</sub>$  on an alkaline earth oxide (BaO, CaO), mediated by precious metals (Pt, Rh), in the form of nitrate or nitrite. This is followed by injection of a rich pulse of a shorter duration to reduce the stored  $NO<sub>x</sub>$  to a mixture of  $N<sub>2</sub>$ ,  $N<sub>2</sub>$ O, and NH<sub>3</sub>. This cycle is repeated continuously to achieve exhaust cycle-averaged  $NO<sub>x</sub>$  conversion greater than 90%.

The development of predictive lean  $NO<sub>x</sub>$  trap (LNT) models is essential to understand and improve the NSR technology. LNT models of varying complexity have been developed by a number of research groups as highlighted in the surveys of Guthenke et al. [\[1\]](#page-14-0) and Roy and Baiker [\[2\]](#page-14-0). Global reaction models provide an assessment of spatio-temporal features while keeping the chemistry manageable. A global LNT model that was developed by Tuttlies et al. [\[3\]](#page-14-0) was based on a shrinking core concept at the particle level to describe the chemistry. A similar model was used by Olsson et al. [\[4\]](#page-14-0) to explain  $NO<sub>x</sub>$  storage. Bhatia et al. [\[5\]](#page-14-0) assumed fast and slow storage sites in their global kinetic model for the LNT. The model was able to predict the spatio-temporal concentration profiles of both reactants and products. Koci et al. [\[6\]](#page-14-0) incorporated the effect of ceria and barium in their comprehensive global model of NSR. Recently, Bhatia et al. [\[7\]](#page-14-0) used global kinetics and crystallite-scale details to explain the effect of Pt dispersion on the rate limiting regime and selectivity for  $Pt/BaO/Al_2O_3$  catalyst.

A comprehensive understanding of NSR mechanism and kinetics ultimately requires a microkinetic description. Xu et al. [\[8,9\]](#page-14-0) developed a microkinetic model for steady state NO reduction by  $H<sub>2</sub>$  in the presence of O<sub>2</sub>. Larson et al. [\[10\]](#page-14-0) also simulated the steady state NO<sub>x</sub> reduction by H<sub>2</sub> and CO in the presence of O<sub>2</sub> using microkinetic formulation. Lindholm et al. [\[11\]](#page-14-0) developed a detailed microkinetic model for the NSR in the presence of  $H_2O$  and  $CO_2$ . These more complex models are continuing to evolve as more detailed transient experiments are conducted to justify their development.

Temporal analysis of products (TAP) was developed by Gleaves and coworkers [\[12\],](#page-14-0) and has been effective in understanding the kinetics of gas–solid catalytic reaction systems. TAP operates under the Knudsen diffusion transport regime, hence eliminating the complexity of mass transfer limitations. Since NSR is an inherently transient process, that makes TAP a particularly suitable tool for mechanistic and kinetic investigations of NSR. In our group, we have employed TAP to study NSR chemistry on  $Pt/Al_2O_3$  and Pt/  $BaO/Al<sub>2</sub>O<sub>3</sub>$  catalysts using NO pulse and NO–H<sub>2</sub> pump–probe experiments [\[13–15\]](#page-14-0). TAP was also coupled with isotopic labeling to gain understanding of spillover mechanisms and distribution of stored  $NO<sub>x</sub>$  in the barium phase [\[16\].](#page-14-0) Recently, TAP was used to estimate the effective diffusivity of stored  $NO<sub>x</sub>$  in the barium phase of  $Pt/BaO/Al_2O_3$  [\[17\].](#page-14-0)



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# Nomenclature



TAP has been used by several researchers to estimate gas diffusivity in porous materials [\[18–23\],](#page-14-0) active site concentration [\[24,25\]](#page-14-0), adsorption and desorption rate constants [\[26\]](#page-14-0), and reaction rate constants using first order kinetics [\[27\].](#page-14-0) Zou et al. [\[28,29\]](#page-14-0) studied pulse gas transport in TAP system. Gleaves and coworkers [\[12\]](#page-14-0) were the first to use exit flux moment analysis to estimate kinetic parameters of heterogeneous reaction systems. Later, a thin-zone reactor model [\[30\]](#page-14-0) was developed to simplify the parameter estimation. In the limit, reactant conversion using a single Pt particle was studied in the reactor filled with inert particles [\[31\].](#page-14-0) Schuurman and Gleaves [\[32\]](#page-14-0) compared the steady state and unsteady state kinetics of n-butane oxidation over VPO using the TAP-2 reactor. Constales et al. [\[33\]](#page-14-0) used the global transfer matrix formulation for single pulse experiments. Feres et al. [\[34\]](#page-14-0) used probability theory to explain the high reactivity of a single Pt particle in the reactor. Kondratenko and Pérez-Ramírez [\[35\]](#page-14-0) modeled N2O decomposition over steam activated Fe–silicate. Hong et al. [\[26\]](#page-14-0) used COMSOL™ to study  $H_2$  and CO adsorption on silica supported cobalt catalysts. Soick et al. [\[36\]](#page-14-0) modeled methane oxidation using 11 species and 14 rate parameters for MgO supported Pt. Balcaen et al. [\[37\]](#page-14-0) used transient multivariable non-linear regression to study oxidative dehydrogenation of propane over vanadia based catalyst.

In spite of the progress made to model gas–solid reactions in the TAP system, there is still a lack of studies focused on quantitative estimation of kinetic parameters for a non-linear reaction system on a multifunctional catalyst. The simplification of a nonlinear system to linear system has a limited validity. The current study is the step to elucidate the effect of various operating parameters for a non-linear reaction system consisting of several surface and gas phase species. Here, we use multiple pulse experiments for NO pulsing and NO–H2 pump–probe to estimate the key kinetic parameters for  $NO-H_2$  system on  $Pt/Al_2O_3$  catalyst. Further, effect of various parameters like  $NO-H<sub>2</sub>$  delay time,  $H<sub>2</sub>/NO$  ratio, and temperature on the product selectivity is discussed. The implications of the findings on the NSR reaction system are discussed.

#### 2. Experimental

#### 2.1. Experimental description

A Generation-1 TAP reactor system was used following the methodologies described in more detail elsewhere [\[13\]](#page-14-0). Experiments were carried out on 1.52 wt.%  $Pt/Al_2O_3$  powder catalyst having 30.4% Pt dispersion,  $1.14 \text{ m}^2$  exposed Pt surface area per gram of catalyst, and 3.7 nm average Pt particle diameter. The catalyst sample was provided by BASF Catalyst LLC (Iselin, NJ). About 16.2 mg of catalyst, with a particle size of  $50-100 \mu m$  and having about 2.3  $\times$  10<sup>17</sup> exposed Pt sites, was sandwiched between two sections of inert quartz particle zones in a cylindrical reactor of 42 mm long and 5.3 mm in diameter. The quartz particles were  $125-150$   $\mu$ m in diameter. The approximate thickness of the catalyst zone is 1 mm, which is sufficiently thin to ensure isothermal operation.

In a typical experiment, catalyst was first oxidized by flowing oxygen over the reactor bed at  $350^{\circ}$ C for 10 min. The oxidation was followed by reduction by multiple pulses of  $H_2$  (typically,

50–100 pulses of H<sub>2</sub> of pulse size  $5 \times 10^{15}$  molecules/pulse) enough to reduce only the exposed Pt and not the alumina as well. The temperature of the catalyst zone was then decreased to the reaction temperature (150–350  $\degree$ C). The components Ar or NO and/or  $H_2$  were injected into the reactor using fast pulsing valves. Effluent species of Ar ( $m/e = 40$ ), N<sub>2</sub> ( $m/e = 28$ ), NO ( $m/e = 30$ ), N<sub>2</sub>O  $(m/e = 44)$ , NH<sub>3</sub> ( $m/e = 16$ ), H<sub>2</sub>O ( $m/e = 18$ ), and H<sub>2</sub> ( $m/e = 2$ ) were monitored with a calibrated UTI 100C quadrupole mass spectrometer. The NH<sub>3</sub> signal was measured at  $m/e = 16$  to avoid overlap with H<sub>2</sub>O at  $m/e = 17$ . The gaseous species signals were calibrated with several hundred pulses on the quartz bed reactor. The inlet pulse sizes were calculated by measuring the pressure drop in an isochoric bulb for each gaseous species. The pulse size was correlated with the mass spectrometer signal using a calibration number. The calibration numbers for  $NH<sub>3</sub>$  and  $H<sub>2</sub>O$  were estimated by N and O balances, respectively, for selected experiment on Pt/  $Al_2O_3$  catalyst and used throughout the study. Typically, the mass balances with this approach closed within 10%. Two principal experimental types were carried out, and are described below.

# 2.2. NO Pulsing experimental results

NO pulsing experiments involve exposure of multiple pulses of NO on pre-reduced catalyst with fixed spacing time  $(\tau_s)$ . The typical size of NO pulse was 1.0  $\times$  10<sup>16</sup> molecules/pulse. The spacing time between consecutive NO pulses was kept at 4.0 s.

These experiments reveal the progressive uptake of NO, O, and N species on Pt and their effect on the evolution of  $N_2$  and  $N_2O$ . NO uptake over Pt/Al<sub>2</sub>O<sub>3</sub> in the temperature range 150–350 °C has been discussed by Kumar et al. [\[15\]](#page-14-0) and Medhekar et al. [\[14\]](#page-14-0) in detail.  $N_2$  is produced during initial NO pulses at high temperature (350 $\degree$ C) and its formation declines with prolonged NO exposure due to oxygen accumulation on the Pt surface. As the surface begins to be covered by oxygen, unreacted NO starts to appear in the effluent and breakthrough of NO is evident. The  $N_2O$  production goes through maxima at the point of NO breakthrough. At a lower temperature, NO uptake decreases, the formation of  $N_2$  decreases and the NO breakthrough occurs at lower pulse number. At 150 °C, the formation of  $N_2$  and  $N_2O$  is negligible due to the kinetic limited N–O bond scission. During the NO uptake experiments for the temperature range considered there is no production of  $O_2$  and NO<sub>2</sub>. Below ca. 400 °C the  $O_2$  desorption is kinetically limited while gaseous  $NO<sub>2</sub>$  formation is thermodynamically limited under TAP conditions. This is in contrast to atmospheric bench scale NO oxidation studies on  $Pt/Al_2O_3$  in which  $NO<sub>2</sub>$  is the major product. If  $NO<sub>2</sub>$  is introduced in the TAP reactor system, it readily decomposes to form NO,  $N_2$ , and  $O_2$  [\[13,15,16\]](#page-14-0).

#### 2.3. NO– $H_2$  pump–probe experimental results

In NO–H<sub>2</sub> pump–probe experiments, pulses of NO and H<sub>2</sub> were sequentially fed to the catalyst. The duration between successive NO and H<sub>2</sub> pulses is the delay time ( $\tau_d$ ) while the time duration between consecutive H<sub>2</sub> and NO pulses is the spacing time ( $\tau_s$ ). In this study,  $\tau_s$  was fixed at 4.0 s while  $\tau_d$  was varied between 0.0 s and 4.0 s.

The NO–H2 pump–probe experiment reveals the transient reaction between NO and  $H_2$  which result in the product distribution that depends on the temperature,  $H_2/NO$  ratio, and the delay time between NO and  $H<sub>2</sub>$  pulses. Higher temperatures favor NO decomposition to produce surface N and O, while at lower temperatures, N–O bond scission is kinetically limited. At higher temperatures (350 °C), surface N combines to from  $N_2$ .  $H_2$  present in the feed is effective in scavenging surface oxygen producing  $H_2O$ . In contrast, at lower temperatures (150 °C),  $H_2$  reacts with surface NO forming  $NH<sub>3</sub>$  as a major product. Hence,  $N<sub>2</sub>$  has maximum selectivity at higher temperatures while  $NH<sub>3</sub>$  has higher selectivity at lower temperatures. Moreover, the N–O bond scission is enhanced by presence of the H adatoms [\[15\]](#page-14-0), leading to  $N_2$  and  $NH_3$ . The  $H_2$ / NO ratio is a critical parameter affecting the  $NH<sub>3</sub>$  and  $N<sub>2</sub>$  selectivity. A higher  $H_2/NO$  feed ratio (>2) leads to higher H surface concentration in proximity to NO or N molecules, resulting in  $NH<sub>3</sub>$  formation. For a smaller ratio of  $H_2/NO$  (<1.5),  $H_2$  scavenges only the O deposited by NO decomposition forming  $H_2O$ . In addition to the temperature and  $H_2/NO$  feed ratio, the delay time between NO and  $H_2$  has a crucial effect on the  $N_2$  and  $NH_3$  selectivity. A longer delay between NO and  $H_2$  facilitates NO decomposition to occur on Pt to form  $N_2$  and surface O that can be scavenged by  $H_2$  pulse. A shorter delay enables an overlap of NO and  $H_2$  on surface that leads to surface interaction of N or NO with H species to form  $NH<sub>3</sub>$ . Hence, ammonia selectivity increases with increase in  $H_2/NO$  ratio and decrease in delay times and temperature.

#### 3. Model development

### 3.1. Reactor model

The TAP reactor is a three-zone fixed-bed reactor (Fig. 1) operated under high vacuum conditions. The predominant gas transport process is Knudsen diffusion. This section summarizes the generalized model of the reactor. The assumptions for the model are

- i. Convection effects are negligible in the reactor and Knudsen diffusion is the predominant transport process.
- ii. Gas phase reactions do not take place.
- iii. Constant physical properties.
- iv. No radial temperature and concentration gradients in the reactor.
- v. No heat or mass transfer limitations.
- vi. Negligible diffusional gradients (or pore diffusional limitations) in the catalyst powder. In otherwords, all of the exposed Pt sites are equally accessible. For the catalyst particle size used this is considered a reasonable approximation.
- vii. Uniform catalyst temperatures.

We used three-zone model for the analysis. A simpler and numerically inexpensive thin-zone model can also be used. There was evidence for concentration gradients along the axis of the reactor for NO adsorption in the catalyst zone. Therefore, we used a more general three-zone model.

The temperature of the reactor is measured in the catalyst zone. The reactor heater is wrapped around central section of the reactor, so the heating is not uniform along the reactor length. Moreover, there are heat losses from the inlet and exit end of the reactor. These factors lead to non-uniformity in the temperature along the reactor. The non-uniformity in temperature for the similar system was reported by Schuurman et al. [\[38\].](#page-14-0) In this study, we assumed the temperature along the reactor is same as the catalyst zone temperature; we will justify the uniform temperature



Fig. 1. Schematic of three-zone powder reactor. Catalyst is sandwiched between two sections of inert zones.

assumption in Section [3.2](#page-4-0). The fluid phase species balances in two inert zones, zone 1 and zone 3, respectively, are given by

$$
\frac{\partial \widehat{C}_k}{\partial t'} = D_{l,k} \frac{\partial^2 \widehat{C}_k}{\partial x_1^2}; \quad 0 < x_1 < L_1,\tag{1}
$$

and

$$
\frac{\partial \widehat{C}_k}{\partial t'} = D_{I,k} \frac{\partial^2 \widehat{C}_k}{\partial x_3^2}; \quad 0 < x_3 < L_3. \tag{2}
$$

The fluid phase species balance of the sandwiched catalyst zone is given by

$$
\frac{\partial \widehat{C}_k}{\partial t'} = D_{c,k} \frac{\partial^2 \widehat{C}_k}{\partial x_2^2} + \mathcal{R}_{g,k} \left( \frac{1 - \varepsilon_c}{\varepsilon_c} \right); \quad 0 < x_2 < L_2,\tag{3}
$$

and the surface species balance in the catalyst zone is given by

$$
\frac{\partial \theta_k}{\partial t'} = \frac{\mathcal{R}_{s,k}}{C_{Pt}},\tag{4}
$$

with the constraint

$$
\sum_{k} \theta_k + \theta_{\nu} = 1, \tag{5}
$$

where the subscript  $\nu$  denotes the vacant sites and the summation is over all surface species.  $\widehat{C}_k$  is the concentration (mol/m<sup>3</sup>) of gas species k in the void space of inert beads or catalyst particles and  $\theta_k$  is the fractional surface coverage of adsorbed species k.  $x_1, x_2,$ and  $x_3$  are length scales (m) in zones 1, 2, and 3, respectively and t' is the time scale (s).  $L_1$ ,  $L_2$ , and  $L_3$  are lengths (m) of zones 1, 2, and 3 respectively.  $D_{I,k}$  and  $D_{C,k}$  are the diffusivities (m<sup>2</sup>/s) of gaseous species k in inert and catalyst zones, and are assumed to be identical. The catalyst zone is very thin  $(\sim 1 \text{ mm})$  compared to the reactor length (L),  $\varepsilon$ <sub>*i*</sub> and  $\varepsilon$ <sub>*c*</sub> are the bed porosity in the inert zones and catalyst zone, respectively. The measured values of porosity are similar. The rate of formation of gas phase species  $k$ ,  $\mathcal{R}_{g,k}$ , and the rate of formation of surface species  $k$ ,  $\mathcal{R}_{s,k}$  are expressed in mol/m<sup>3</sup> catalyst/s.  $C_{Pt}$  is the exposed Pt surface concentration in mol Pt/ $m<sup>3</sup>$  catalyst. The reactor parameters used in this study is listed in Table 1.

Reactant or inert gases are pulsed into the reactor at  $x_1 = 0$ . The pulse width is 200-500 µs; this width is very small in comparison to diffusion time scale ( $\sim$ 0.1 s), and hence the shape of inlet pulse (e.g. modified Gaussian or positive sinusoidal) does not affect the transient effluent shape. Even the pulse width interval does not change the effluent profiles as long as it is much smaller than the diffusion time scale. For the purpose of this study, we assume a positive sinusoidal pulse shape that is continuous and differentiable in the time interval of pulse width. A non-differential pulse shape like triangular or square pulse shape would cause steep flux gradients at the reactor inlet and would require smaller time steps for the numerical integration. In order to decrease computation time,



Properties and reactor parameters used in this study.



those pulse shapes are avoided. The boundary condition at the inlet of reactor is defined as

$$
-A\varepsilon_{I}D_{I,k}\frac{\partial\widehat{C}_{k}}{\partial x_{1}}\bigg|_{x_{1}=0} = N_{k}\delta_{k}(t'),
$$
\n(6)

where  $\delta_k$  is the positive sinusoidal pulse. A is the cross sectional area of the reactor and  $N_k$  is the pulse size (mol/pulse) for species k. The outlet of the reactor is exposed to high vacuum of  $10^{-8}$  torr. So, the concentration of gaseous species at the reactor exit is essentially zero, i.e.,

$$
\widehat{C}_k|_{x_3=L_3}=0.\tag{7}
$$

In addition to the above inlet and outlet boundary conditions, the concentration and flux continuity are defined at the zones 1– 2 interfacial boundary

$$
\hat{C}_k|_{x_1 = L_1^-} = \hat{C}_k|_{x_2 = 0^+},\tag{8}
$$

and

$$
\varepsilon_{i}D_{l,k}\frac{\partial\widehat{C}_{k}}{\partial x_{1}}\bigg|_{x_{1}=L_{1}^{-}}=\varepsilon_{C}D_{C,k}\frac{\partial\widehat{C}_{k}}{\partial x_{2}}\bigg|_{x_{2}=0^{+}}.
$$
\n(9)

Similarly, the gaseous concentration and flux continuity at the zones 2–3 interface are defined as

$$
\hat{C}_k|_{x_2 = L_2^-} = \hat{C}_k|_{x_3 = 0^+},\tag{10}
$$

and

$$
\varepsilon_{c}D_{c,k}\frac{\partial\widehat{C}_{k}}{\partial x}\Big|_{x_{2}=L_{2}^{-}}=\varepsilon_{l}D_{l,k}\frac{\partial\widehat{C}_{k}}{\partial x}\Big|_{x_{3}=0^{+}}.\tag{11}
$$

The concentration of gas and surface species in the reactor bed at the start of the pulse is equal to gas or surface concentration at the end of the previous cycle. The cycle time is  $\tau_s$  for a pulsing experiment and  $\tau_d + \tau_s$  for pump–probe experiment. As before,  $\tau_s$ is the spacing time and  $\tau_d$  is the delay time. The initial conditions for the gas concentration and surface concentration are

$$
\left.\hat{C}_k\right|_{\substack{t'=0 \text{pulse}=p+1}} = \hat{C}_k|_{t' = \tau_d + \tau_s},\tag{12}
$$

and

$$
\theta_k \Big|_{\substack{t'=0 \text{pulse}=p+1}} = \theta_k \Big|_{\substack{t' = \tau_d + \tau_s:\\ \text{pulse}=p}} \tag{13}
$$

In addition to the calculation of gaseous and surface concentration profiles, the exit flux (mol/s) at the reactor outlet is calculated as

$$
F_k = -A\varepsilon_l D_{l,k} \frac{\partial \widehat{C}_k}{\partial x_3} \bigg|_{z_3 = L_3}.
$$
\n(14)

The above equations are non-dimensionalized using the following quantities

$$
C_o = \frac{N_o}{A\epsilon_l L},\tag{15}
$$

$$
z_1 = \frac{x_1}{L},\tag{16}
$$

$$
z_2 = \frac{L_1 + x_2}{L},\tag{17}
$$

$$
z_3 = \frac{L_1 + L_2 + x_3}{L},\tag{18}
$$

$$
t = \frac{t'}{L^2/D_o},\tag{19}
$$

<span id="page-4-0"></span>and

$$
C_k = \frac{\widehat{C}_k}{C_o}.
$$
 (20)

Here,  $L$  is the length (m) of the reactor,  $C<sub>o</sub>$  is the reference concentration (mol/m $^3$ ),  $N_o$  is reference pulse size (mol/pulse) and  $D_o$  is the reference diffusivity ( $m^2/s$ ).  $z_1$ ,  $z_2$ , and  $z_3$  are dimensionless length scales for zones 1, 2, and 3, respectively, and  $t$  is the dimensionless time scale. Using the above expressions, the non-dimensionalized equations for zones 1, 2, and 3 are given by

Zone 1:

$$
\frac{\partial C_k}{\partial t} = \frac{D_{I,k}}{D_o} \frac{\partial^2 C_k}{\partial_{x_1}^2}; \quad 0 < z_1 < \frac{L_1}{L},\tag{21}
$$

Zone 2:

$$
\frac{\partial C_k}{\partial t} = \frac{D_{C,k}}{D_o} \frac{\partial^2 C_k}{\partial_{x_2}^2} + \mathscr{R}_{g,k} \frac{L^2}{C_o D_o} \left( \frac{1 - \varepsilon_C}{\varepsilon_C} \right); \quad \frac{L_1}{L} < z_1 < \frac{L_1 + L_2}{L},\tag{22}
$$

$$
\frac{\partial \theta_k}{\partial t} = \frac{L^2 \mathcal{R}_{s,k}}{D_o C_{Pt}},\tag{23}
$$

with the constraint

$$
\sum_{k} \theta_k + \theta_{\nu} = 1. \tag{24}
$$

Zone 3:

$$
\frac{\partial C_k}{\partial t} = \frac{D_{l,k}}{D_o} \frac{\partial^2 C_k}{\partial_{x_3}^2} ; \frac{L_1 + L_2}{L} < z_3 < 1.
$$
 (25)

Boundary conditions are as follows

Reactor inlet:

$$
\left. \frac{\partial C_k}{\partial z_1} \right|_{z_1=0} = -\frac{D_o}{D_{I,k}} \frac{N_k}{N_o} \delta_k(t). \tag{26}
$$

Here,  $\delta_k(t)$  is defined for an inlet pulse as

$$
\delta_k(t) = \frac{1}{2\Delta t_{Value}} \left( 1 + \sin\left(\frac{\pi}{\Delta t_{Value}} \left(t - \frac{\Delta t_{Value}}{2}\right)\right) \right), \quad 0 \leq t
$$
\n
$$
\leq 2\Delta t_{Value}, \tag{27}
$$

or the finite Gaussian pulse

$$
\delta_k(t) = \frac{1.3027}{\Delta t_{Value}} e^{-\frac{4(t-\Delta t_{Value})^2}{t(2\Delta t_{Value}-t)}}, \quad 0 \leq t \leq 2\Delta t_{Value},
$$
\n(28)

where  $\Delta t_{Value}$  is the half of dimensionless opening time of pulse valve. For this simulation, we used a positive sinusoidal response.

Zones 1–2 interface:

$$
C_k\big|_{z_1 = \frac{L^-}{L}} = C_k\big|_{z_2 = \frac{L^+}{L}},\tag{29}
$$

and

$$
\left. \frac{\partial C_k}{\partial z_1} \right|_{z_1 = \frac{l_1^-}{l}} = \frac{\varepsilon_c D_{C,k}}{\varepsilon_l D_{l,k}} \frac{\partial C_k}{\partial z_2} \bigg|_{z_2 = \frac{l_1^+}{l}}.
$$
\n(30)

Zones 2–3 interface:

$$
C_k\big|_{z_2 = \frac{L_1 + L_2^-}{L}} = C_k\big|_{z_3 = \frac{L_1 + L_2^+}{L}},\tag{31}
$$

$$
\left. \frac{\partial C_k}{\partial z_2} \right|_{z_2 = \frac{L_1 + L_2}{L}} = \frac{\varepsilon_1 D_{I,k}}{\varepsilon_C D_{C,k}} \frac{\partial C_k}{\partial z_3} \bigg|_{z_3 = \frac{L_1 + L_2^+}{L}}.
$$
\n(32)

Reactor exit:

$$
C_k|_{z_3=1}=0.
$$
\n(33)

Initial conditions:

$$
C_k \vert_{\substack{t=0 \text{pulse}=P+1}} = C_k \vert_{\substack{t=\frac{\tau_d+\tau_s}{L^2/D_0}, \text{pulse}=P}} \tag{34}
$$

and

$$
\theta_k \Big|_{\substack{t=0 \text{pulse}=p+1}} = \theta_k \Big|_{\substack{t=\frac{\tau_d+\tau_s}{L^2/D_o}.}} \tag{35}
$$

Finally, the exit flux of species  $k$ ,  $F_k$ , can be calculated using

$$
F_k = -A\varepsilon_l D_{l,k} \frac{C_o}{L} \frac{\partial C_k}{\partial z_3}\bigg|_{z_3=1}.
$$
\n(36)

#### 3.2. Transport regime

The predominant transport process in the reactor is Knudsen diffusion. This implies that the mean free path of the gas molecules is greater than the characteristic catalyst particle diameter in the reactor. This type of transport regime rules out the possibility of gas–gas interactions inside the reactor, so the primary interactions taking place in the reactor are of the solid–gas variety. The pulses of Ar with inlet pulses of size  $2 \times 10^{15}$ – $2 \times 10^{16}$  molecules/pulse were fed to the three-zone reactor at  $350^{\circ}$ C and the exit flux was recorded. The area normalized flux responses were identical for each pulse size. The similar identical area normalized profiles were also obtained for N<sub>2</sub> at 350 °C and Ar at 250 °C and 150 °C. This means that the pulse size (up to  $2 \times 10^{16}$  molecules/pulse) does not affect the area normalized pulse shape, which is a characteristic of Knudsen diffusion transport. Further, the peak time  $(t'_{p,k})$ and area normalized peak flux  $(J_{p,k})$  for each case results in  $J_{p,k}\times t_{p,k}^\prime$  = 0.31. Hence, assumption of Knudsen diffusion for the above TAP model development is justified.

The temperature and molecular weight is related to Knudsen diffusivity as

$$
D_k = \phi \sqrt{\frac{T}{M_{w,k}}},\tag{37}
$$

where  $\phi$  is an empirical parameter that depends on bed porosity, tortuosity, particle size, etc. Gleaves et al. [\[39\]](#page-14-0) derived an expression for the peak time  $(t'_{p,k})$  dependence on the reactor length  $(L)$ and diffusivity  $(D_k)$  as

$$
t'_{p,k} = \frac{L^2}{6D_k}.\tag{38}
$$

From the above two equations, peak time correlates to  $\frac{M_{w,k}}{T}$  $\mathcal{L}_{\mathcal{A}}$ as follows

$$
t'_{p,k} = \frac{L^2}{6\phi} \sqrt{\frac{M_{w,k}}{T}}.
$$
\n(39)



Fig. 2. Dependence of peak time on molecular weight and temperature for diffusivity estimation.



Fig. 3. Exit flux profiles for inert gases (Ar and  $N_2$ ). The symbols are the experimental data and lines are simulation results throughout this paper.

The value of  $\phi$  from [Fig. 2](#page-4-0) is estimated as 7.77  $\times$  10<sup>–4</sup> m<sup>2</sup>/s (g/mol/  $K$ <sup>0.5</sup>. This value is used to fit the transient diffusion curves for Ar at 150–350 °C and N<sub>2</sub> at 350 °C. The experimental transient curves fit very well with the simulated curves with the value of  $\phi$  (Fig. 3). This value is used throughout this study for diffusivity estimation.

Although the temperature in the reactor is not uniform, the estimated diffusivity using a single temperature (measured in catalyst zone) is able to capture the exit flux profiles of inert species like  $N_2$ and Ar (Fig. 3). For the case of reacting species in the Knudsen regime, the non-uniform temperature does not change the transport behavior. The non-uniform temperature can result in non-uniform gas–solid as well as gas–gas phase reaction rates. However, the quartz is an inert at the temperatures used, therefore, the gas–solid reactions are absent in the quartz zones. The gas–gas phase reactions are absent due to Knudsen transport conditions. The only reactions that occur are gas–solid reactions in the catalyst zone. Moreover, the temperature is measured in catalyst zone; therefore, the reaction rates are not affected by non-uniform temperature effects.

In addition to the uniform temperature assumption, we assumed constant diffusivity in the catalyst zone and inert zone. The nominal diameter of inert quartz particle is  $125-150$   $\mu$ m while the catalyst particle size was less than  $100 \mu m$ . The difference in the particle diameter results in difference in diffusivity in the inert zone and catalyst zone. The thickness of catalyst zone  $(\sim 1 \text{ mm})$  is very small in comparison to reactor length (42 mm). A single uniform diffusivity in the reactor can explain the exit flux behavior for  $N_2$  and Ar pulses in temperature range 150–350 °C (Fig. 3). Therefore, two different diffusivities, each for inert and catalyst zones, are not needed.

# 4. Modeling results and discussion

# 4.1. NO decomposition

NO decomposition experiments reveal the progressive uptake of NO, O, and N species on Pt and their effect on the evolution of  $N_2$  and  $N_2$ O. NO decomposition at 250 and 350 °C involves stoichiometric formation of  $N_2$  during the initial pulses followed by breakthrough of NO. The production of  $N_2O$  starts at zero, achieves a maximum during NO breakthrough, and then slowly decreases with pulse number (Fig. 4).

As described earlier,  $NO<sub>2</sub>$  and  $O<sub>2</sub>$  are not feasible products under TAP pressure and temperature conditions. The simplest microki-



Fig. 4. Effluent integral profiles for NO pulsing at 350  $\degree$ C, symbols are experimental data while lines are calculated using kinetic parameters from Xu et al. [\[9\]](#page-14-0).

netic mechanism for  $N_2$  and  $N_2O$  formation involves reversible NO adsorption, NO decomposition,  $N_2$  formation, and  $N_2O$  formation. A four step mechanism with associated reaction rate parameters from Xu et al. [\[9\]](#page-14-0) was used as a starting point. This was a subset of the larger microkinetic model for anaerobic NO reduction by  $H_2$  on alumina-supported Pt. The model captured most experimental trends under steady state anaerobic conditions, particularly under rich conditions ( $H_2/NO \ge 1$ ). In the NO decomposition submodel the formation of  $N_2$  is considered as an irreversible reaction, while the reversible adsorption of NO is non-activated. The microkinetic mechanism for NO decomposition and the corresponding rate parameters used by Xu et al. [\[9\]](#page-14-0) are provided in Table 2.

The integral profiles of  $N_2$ , NO, and  $N_2$ O predicted using the ori-ginal kinetic parameters from Xu et al. [\[9\]](#page-14-0) at 350  $\degree$ C are shown in Fig. 4. During the first 15 pulses, the  $N_2$  production is constant and unreacted NO in the effluent is negligible. The intersection of the NO and  $N_2$  profiles occurs at the 20th pulse and the  $N_2O$  production is negligible over all the pulses. On the other hand, the experimental data have somewhat different trends. The data show a finite production of  $N_2O$ , gradual decrease in  $N_2$  starting with the first pulse, and a gradual increase in effluent NO flux. Obviously, some tuning of the kinetic parameters and potential upgrade in the mechanism is needed to improve the simulation of the integral pulse data. A comparison of individual species flux response also showed the need for parameter tuning.

The next step involved a more rigorous fit of the model predictions to the transient species fluxes. A multi-response non-linear least squares algorithm was devised based on minimizing the least square sum of the objective function  $(e)$  between experimentally measured and simulated values of individual fluxes (mol/s). The objective function e

an c	
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Four step mechanism and rate constants for NO decomposition adopted from Xu et al. [9].



<span id="page-6-0"></span>
$$
e = \sum_{\text{Pulse number Temperature}} \sum_{\text{Gas species Time}} \sum_{\text{Time}} w(F^{\text{Exp}} - F^{\text{Sim}})^2. \tag{40}
$$

The summation was over selected pulses of NO,  $N_2$ , and  $N_2O$  for initial 30 pulses; two temperatures 250  $\degree$ C and 350  $\degree$ C; three gas phase species and 390 data points per pulse transient. Though we collected experimental data with 1 ms interval, but for the simulation we use data points with 10 ms interval since it adequately capture the transient response of all species. In the above equation, w is the weight,  $F^{Exp}$  is the experimentally observed exit flux (mol/s), and  $F^{Sim}$  is the exit flux (mol/s) estimated by simulation. Specifically, the NO decomposition rate parameters were estimated for the individual transient pulses of NO,  $N_2$ , and  $N_2O$  at 250 and 350 °C using initial guess values from Xu et al. [\[9\]](#page-14-0). Another choice of the objective function may be defined between experimentally measured and simulated values of integral fluxes (mol/pulse). That would not be a good choice because the simulated transient fluxes (mol/ s) may not necessary match to the experimentally measured transient fluxes though the integral fluxes (mol/pulse) are in good match. It is noted that pore diffusion limitations may have been present for some conditions, but the extent is likely not to be significant given the  $50-100 \mu m$  particle size.

We used DLSODE to solve the system of PDE and ODRPACK 2.01 to optimize the rate parameters. The weights were selected based on an inspection of the typical responses. The rationale was as follows. A typical effluent flux profile during the 0–4 s period showed NO and  $N_2$  as the major effluent species while  $N_2O$  fluxes were much lower (Fig. 5). The fluxes of NO and  $N_2$  were highest in 0– 1 s range and decreased to the baseline level afterwards (Fig. 6). For this reason, a base weight 1 was assigned for the NO and  $N<sub>2</sub>$ fluxes in the 0–1 s interval and assigned a lower value of 0.2 afterwards. The  $N<sub>2</sub>O$  flux was generally an order of magnitude smaller than the NO and  $N_2$ . Hence, a larger weight of 8 was assigned for N2O in the 0–1 s interval, 2 in 1–2 s interval, and 1 in 2–4 s interval. Similar logic was used for weight assignment for NO pulsing experiment at 250 $\degree$ C also. The weights are tabulated in [Table 3](#page-7-0). There are more formal ways for optimization for these transient multiresponse curves like Bayesian multi-response minimization that would potentially better fit the data and does not require weight assignment.

The initial kinetic parameters optimization for NO decomposition experiment was tried using steps D1–D4. The optimized parameters were not able to simultaneously explain all the features of NO decomposition experiment listed above. Hence, a need



Fig. 5. NO decomposition on Pt/Al<sub>2</sub>O<sub>3</sub> at 350 °C (a) and 250 °C (b). Symbols are experimental data while solid lines are simulation results throughout the paper.



Fig. 6. Experimental and simulation results for individual transient pulse profiles for NO, N<sub>2</sub>, and N<sub>2</sub>O at 350 °C.

Tab

<span id="page-7-0"></span>Table 3 Weights of individual flux responses for various temperatures, species, and transient time.

	250 °C			350 $\degree$ C		
Time $(s)$	N <sub>2</sub>	N <sub>O</sub>	N <sub>2</sub> O	Nэ	NO	N <sub>2</sub> O
$0.0 - 1.0$			x			8
$1.0 - 2.0$				0.2	0.2	
$2.0 - 3.9$				0.2	0.2	

to change the reaction mechanism was realized. We added an additional reaction for  $N_2O$  formation; i.e. disproportionation of NO to form  $N_2O$  (Reaction D5). The resulting five-step reaction mechanism (D1–D5) is summarized in Table 4. The corresponding rate expressions are provided in Table 5. Here,  $R_{if}$  and  $R_{if}$  are the forward and reverse rates for reaction *i*, in mol/m<sup>3</sup>/s.

In the parameter estimation the forward activation energies of steps D2 (adsorbed NO decomposition) and D3 ( $N_2$  formation) were kept the same as those provided by Xu et al. [\[9\].](#page-14-0) We fixed the values of these activation energies in order to avoid convergence problems in the parameter estimation scheme. The NO desorption activation energy was fixed at 100.0 kJ/mol. This value is slightly lower than reported value of 114.5 kJ/mol [\[9\]](#page-14-0). A higher value of NO decomposition activation energy results in a sharper NO breakthrough and a sharper  $N_2$  decrease. In addition, we used one enthalpy balance equation (NO + NH<sub>3</sub>  $\leftrightarrow$  N<sub>2</sub>O + 1<sup>1</sup>/<sub>2</sub> H<sub>2</sub>), that uses rate constants from the NO–H2 reaction kinetics, to estimate the  $N_2O$  adsorption activation energy. This leads to three fixed parameters and 11 optimization parameters for the NO decomposition system.

The rate constant for adsorption of a species k, involved in reaction step  $i$ , is given by

$$
k_i = S_{0,k} \alpha_v \sqrt{\frac{RT}{2\pi M_{w,k}}},\tag{41}
$$

where  $S_{0,k}$  is the sticking coefficient of species k at zero coverage, R is the universal gas constant, T is temperature in K,  $M_{w,k}$  is the molecular weight of species k in kg/mol and  $\alpha_v$  is the active Pt surface area per unit volume of catalyst;  $\alpha_v$  is calculated using

$$
\alpha_{\nu} = \frac{C_{Pt}}{\Gamma},\tag{42}
$$

where  $C_{Pt}$  is the moles of Pt exposed per unit volume of catalyst and  $\Gamma$  is the surface site density in moles Pt per unit area of exposed Pt. The surface site density is an intrinsic parameter for a catalyst. We used a nominal value for  $\Gamma$  available in literature [\[40,41\].](#page-14-0) The rate of adsorption is given by

$$
R_i = k_i \hat{C}_k \theta_v^n, \tag{43}
$$

where  $\widehat{\mathsf{C}}_k$  is the concentration (mol/m<sup>3</sup>) of a gaseous species  $k$  in gas phase and  $\theta_v$  is the fractional coverage of vacant Pt. The order of adsorption,  $n$ , is equal to 2 for second order adsorption and unity for first order adsorption. Here, we defined all the rate of reactions





in terms of surface concentration and gaseous concentration. The Arrhenius rate expression was used for surface reactions, desorption and other adsorption reaction systems in which Eq. (41) was not used;

$$
k_i = A_i e^{-\frac{E_i}{RT}},\tag{44}
$$

 $A_i$  is the pre-exponential factor and  $E_i$  is the activation energy (J/mol) for reaction step i.

The parameter estimation algorithm and upgraded microkinetic model for NO decomposition gave a much improved prediction of the experimental data. [Fig. 5](#page-6-0) shows the integral exit fluxes of  $N_2$ , NO, and N<sub>2</sub>O with pulse number for NO decomposition at 350  $\degree$ C and 250 $\degree$ C. The five-step microkinetic mechanism with estimated rate constants is able to better capture and help us to elucidate the stoichiometric production of  $N_2$  during initial pulse, NO breakthrough, and  $N_2O$  maxima. The additional route for  $N_2O$  formation (D5) by NO disproportionation reaction improves the  $N_2O$  prediction. Our simulation shows that NO disproportionation is the dominating route for  $N_2O$  formation.

The corresponding experimental and model predicted individual pulse responses of NO,  $N_2$ , and  $N_2O$  at 350 °C are shown in [Fig. 6](#page-6-0). The model predicts a  $N_2$  flux profile that decreases with pulse number due to oxygen accumulation, and as a result the unreacted NO in the effluent increases with the pulse number. The model predicts reasonably well the peak flux, peak time, and their variation with the pulse number. Similar results of NO decomposition at 250  $\degree$ C are shown in [Fig. 7](#page-8-0). The least square fit at 350 °C is somewhat better than that at 250 °C. At 250 °C, the spacing time between consecutive NO pulses was kept at 4.0 s and was insufficient for the gaseous NO concentration (or flux) to reach a baseline level. The experimental NO desorption profile has a protracted tail as does the  $N_2$  and  $N_2O$  effluent profiles. The possible reasons may be that desorption of NO from Pt is slow and/or  $Al_2O_3$  is also participating in the adsorption/desorption of NO. In the simulations, we accounted for the pulse not returning to the baseline level by using the instantaneous spatial profile as the initial condition for the next pulse.

In order to improve our understanding of the interaction of NO on Pt, we plot the exit fluxes of NO,  $N_2$ , and  $N_2O$  at 350 °C [\(Fig. 8\)](#page-8-0). The surface coverages of adsorbed NO, N, and O as well as gaseous concentrations of NO,  $N_2$ , and  $N_2O$  are also plotted at the front of catalyst zone for pulses 1, 2, 5, 10, and 20. It is seen that the exit flux of  $N_2$  is highest during the first pulse and decreases monotonically with pulse number. During this span of pulses, the fluxes of

Table 4

Rate parameters for NO decomposition on Pt. Underlined parameters were fixed, italicized parameter was calculated from enthalpy balance and others were optimized.

	Reactions	$A_{i,f}$ (mol/m <sup>3</sup> /s or s <sup>-1 (*)</sup> )	$E_{if}$ (kJ/mol)	$A_{ir}$ (mol/m <sup>3</sup> /s or s <sup>-1 (*)</sup> )	$E_{ir}$ (kJ/mol)
D <sub>1</sub>	$NO + Pt \leftrightarrow NO - Pt$	$(2.38 \pm 0.19) \times 10^{8}$ <sup>(*)</sup>		$(2.45 \pm 0.9) \times 10^{16}$	<u>100.0</u>
D <sub>2</sub>	$NO-Pt + Pt \leftrightarrow N-Pt + O-Pt$	$(8.76 \pm 12.56) \times 10^{13}$	107.8	$(6.42 \pm 10.54) \times 10^{14}$	$104.5 \pm 5.0$
D <sub>3</sub>	$2 N-Pt \rightarrow N_2 + 2Pt$	$(2.71 \pm 1.17) \times 10^{16}$	130.0		
D <sub>4</sub>	$NO-Pt + N-Pt \leftrightarrow N_2O + 2 Pt$	$(3.69 \pm 0.2) \times 10^9$	$81.5 \pm 4.5$	$(1.18 \pm 65.90) \times 10^{3}$ <sup>(*)</sup>	36.2
D <sub>5</sub>	2 NO−Pt → N <sub>2</sub> O + O−Pt + Pt	$(6.45 \pm 0.6) \times 10^{15}$	$152.1 \pm 1.4$		

 $(*)$  denotes the pre-exponential factors ( $A_{i,f}$  and  $A_{i,r}$ ) with unit s<sup>-1</sup>.

<span id="page-8-0"></span>

Fig. 7. Experimental and simulation results for individual transient pulse profiles for NO, N<sub>2</sub>, and N<sub>2</sub>O at 250 °C.



**Fig. 8.** Exit flux, surface coverage ( $\theta$ ), and concentration at front of catalyst zone for NO pulsing at 350 °C with pulse size 1.0 × 10<sup>16</sup>.

NO and  $N_2O$  increase. By the 10th pulse, the NO flux surpasses that of  $N_2$  and continues to grow. The surface coverages of adsorbed N and NO are negligible in comparison to adsorbed O at 350  $\degree$ C. This shows that oxygen is the predominant species present on the Pt surface. The concentration profiles of gaseous species in the catalyst zone are similar to the exit flux profile.

The maximum gaseous concentration of any species for an inlet NO pulse of magnitude 1  $\times$  10<sup>16</sup> molecules/pulse is  $\sim$ 10<sup>22</sup> molecules/m<sup>3</sup>. The total pressure of gases in the catalyst zone can be estimated by

$$
P = \hat{C}RT.\tag{45}
$$

The above equation gives the maximum pressure in the catalyst zone to be about 86 Pa. The mean free path of the gaseous molecules  $(\lambda)$  in the packed bed reactor can be calculated as

$$
\lambda = \frac{RT}{\sqrt{2}\pi d^2 N_A P},\tag{46}
$$

where  $d$  is the diameter of a gaseous molecule and  $N_A$  is Avogadro's number. The mean free path  $(\lambda)$  for a typical gas molecule of molecular diameter 0.3 nm at 350 °C is 250  $\mu$ m. The value of the mean free path is greater than a typical catalyst powder size ( $\sim$ 50 µm) and inert quartz particle ( $\sim$ 125–150 µm). Hence, the assumption <span id="page-9-0"></span>of Knudsen diffusion for pulse size  $10^{16}$  molecules/pulse is still valid. It has been already shown that the area normalized fluxes of  $N_2$ and Ar coincide for the pulse size (<2  $\times$  10<sup>16</sup> molecules/pulse) and the temperatures below 350 °C. Moreover, the product  $(J_{p,k}\times t'_{p,k})$ is equal to 0.31 and the calculated diffusivity explains well the exit flux profiles of inert gases confirms the Knudsen flow regime.

Finally, the current microkinetic mechanism of NO decomposition contains five steps. The mechanism involves three gas phase species (NO,  $N_2$ , and  $N_2O$ ) and two elemental species (N and O). Further, the microkinetic formation of  $N<sub>2</sub>$  is assumed to be irreversible due to the very high adsorption activation energy of  $N_2$ . The thermodynamic constraint involves two gas phase species (NO and  $N_2O$ ) with two elements (N and O). Hence, the NO decomposition reaction system does not require any thermodynamic constraint to be satisfied.

#### 4.2. NO- $H_2$  pump-probe

The NO–H2 pump probe experiments have a rich complexity of surface chemistry with multiple products formed. The experiments resulted in the formation of  $H_2O$  and  $NH_3$  in addition to  $N_2$  and  $N<sub>2</sub>$ O. In this section, we discuss the effect of temperature, delay time and  $NO/H<sub>2</sub>$  ratio on product selectivity. The earlier study of Kumar et al. [\[15\]](#page-14-0) provided related results to those reported here. However, the current data were used for the model simulations.

The three primary overall reactions between NO and  $H<sub>2</sub>$  are

$$
NO + \frac{1}{2}H_2 \rightarrow \frac{1}{2}N_2O + \frac{1}{2}H_2O,
$$
  

$$
NO + H_2 \rightarrow \frac{1}{2}N_2 + H_2O, \text{ and}
$$

$$
NO + 2\frac{1}{2}H_2 \rightarrow NH_3 + H_2O.
$$

The stoichiometric ratio of  $H<sub>2</sub>$  to NO is an important factor determining the formation of  $N_2O$ ,  $N_2$ , or  $NH_3$ . We define the feed ratio  $H_2/NO = 1$  for the production of N<sub>2</sub>. Thus, for  $H_2/NO < 1$ ,  $H_2$  is the limiting reactant and NO is in excess; and for  $H_2/NO > 1$ , NO is the limiting reactant and  $H_2$  is in excess.

The NO–H<sub>2</sub> pump–probe chemistry (P1–P9) involves H<sub>2</sub>O and  $NH<sub>3</sub>$  formation in addition to  $N<sub>2</sub>$  and  $N<sub>2</sub>O$  (Table 6). The corresponding rate expressions are described in Table 7. The kinetic parameters are as it is adopted from Xu et al. [\[9\]](#page-14-0) with the following changes. The reaction (NH<sub>3</sub>–Pt + 3 OH–Pt  $\rightarrow$  N–Pt + 3 H<sub>2</sub>O + 2 Pt) was eliminated because it overpredicts the formation of  $N_2$  and the rate parameters for step P8 were modified to fit the  $N_2$  and  $NH<sub>3</sub>$  integral profile. Here, for the case of NO–H<sub>2</sub> pump–probe chemistry on Pt, no further effort was made to optimize the kinetic parameter estimation.

Further, some of the reaction steps like ammonia formation (P5 and P6) and ammonia consumption (P8 and P9) are not elementary



#### Table 7











steps [\[9\],](#page-14-0) but are steps that lump the sequential addition of H to the adsorbed NO or N and abstraction of H from surface ammonia.

NO– $H_2$  pump–probe involves formation of NH<sub>3</sub> and H<sub>2</sub>O. The delayed appearance of these species suggests that  $NH<sub>3</sub>$  and  $H<sub>2</sub>O$ adsorption/desorption is important. The adsorption of  $H_2O$  and  $NH<sub>3</sub>$  on Pt is not sufficient to explain the breakthrough of NH<sub>3</sub> and  $H<sub>2</sub>O$ , so adsorption on  $Al<sub>2</sub>O<sub>3</sub>$  (steps A1 and A2) had to be taken into account (Table 8). For  $NH<sub>3</sub>$  this is not surprising since it is basic while the  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> has acidic sites, so NH<sub>3</sub> readily adsorbs on Al<sub>2</sub>O<sub>3</sub>. Water as well has an affinity for the  $\gamma$ -Al<sub>2</sub>O<sub>3</sub>. Table 9 provides the rate expressions for  $H_2O$  and  $NH_3$  adsorption on  $Al_2O_3$ . The surface area of Pt is 1.14  $m^2/g$  of catalyst. For this catalyst sample, Pt is dispersed over high surface area  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> and the BET surface area of the catalyst is 100–120 m $^2$ /g of catalyst. This means that the surface area of  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> is approximately 100 times that of the Pt surface area, hence we used surface concentration of  $\text{Al}_2\text{O}_3$  sites as 100 times the concentration of Pt sites. The rate constant for adsorption of  $H_2O$  (step A1) and NH<sub>3</sub> (step A2) on  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> are given by

$$
k_i = S_{0,k}(100\alpha_v)\sqrt{\frac{RT}{2\pi M_{w,k}}}.\tag{47}
$$

Table 9

Rate expressions for  $H_2O$  and  $NH_3$  interactions on  $Al_2O_3$ .

	Reaction	$R_{if}$ (mol/m <sup>3</sup> /s)	$R_{ir}$ (mol/m <sup>3</sup> /s)
A1	$H_2O + Al_2O_3 \leftrightarrow H_2O-Al_2O_3$	$k_{A1,f} \widehat{C}_{H_2O} \theta_{v,Al_2O_3}$	$k_{A1,r}\theta_{H_2O,Al_2O_3}$
A <sub>2</sub>	$NH_3 + Al_2O_3 \leftrightarrow NH_3-Al_2O_3$	$k_{A2,f} \hat{C}_{NH_3} \theta_{\nu,Al_2O_3}$	$k_{A2,r}\theta_{NH_3,Al_2O_3}$



<span id="page-10-0"></span>The adsorption of  $H_2O$  and  $NH_3$  on  $Al_2O_3$  is inactivated. The sticking coefficient of 0.1 is used for water adsorption [\[42\]](#page-14-0) and 1.0 is assumed for NH<sub>3</sub> adsorption. The desorption rate constants were estimated by fitting the  $H_2O$  and NH<sub>3</sub> profile for selected NO– $H_2$  pump– probe experiment that produce both  $H_2O$  and  $NH_3$ . These values of desorption constants are used throughout the study.

The experimental results show a delayed appearance of  $H_2O$ and NH<sub>3</sub> in the effluent. The integral response of H<sub>2</sub>O and NH<sub>3</sub> shows a monotonic increase in their production, eventually approaching a saturation level. In order to match the integral trends of  $H_2O$  and  $NH_3$ , the saturated level values of  $H_2O$  and  $NH<sub>3</sub>$  during the NO–H<sub>2</sub> experiment were used as the inlet pulse sizes of  $H_2O$  and  $NH_3$ , respectively. The Pt chemistry was neglected for estimation of adsorption/desorption response on the  $Al_2O_3$ . The desorption rate constants for  $H_2O$  and  $NH_3$  were estimated by integral responses of H<sub>2</sub>O and NH<sub>3</sub> at 250 °C and 350 °C. The desorption rate constants were plotted against  $\frac{1}{RT}$  to estimate activation energy of desorption and the pre-exponential factor [\(Table 8\)](#page-9-0).

The  $NO-H<sub>2</sub>$  reaction system involves six gas phase species ( $NO$ ,  $N_2$ ,  $N_2$ O,  $NH_3$ ,  $H_2$ O, and  $H_2$ ) and three atomic species (N, O, and H). The formation of  $H_2O$  and  $N_2$  are assumed to be irreversible. As a result, a thermodynamic analysis involves four gaseous species (NO,  $N_2O$ ,  $NH_3$ , and  $H_2$ ) and three elemental species (N, O, and H). This leaves one thermodynamic constraint for reaction  $(NO + NH_3 \leftrightarrow N_2O + 1\frac{1}{2}H_2)$ , which has a heat of reaction of 35.91 kJ/mol. This reaction can be obtained by combining the steps



**Fig. 9.** Effect of H<sub>2</sub>/NO ratio on product selectivity at 250 °C and 0.1 s delay time. (a and b) H<sub>2</sub>/NO = 2.26, NO<sub>In</sub> = 1.09  $\times$  10<sup>16</sup>; (c and d) H<sub>2</sub>/NO = 1.04, NO<sub>In</sub> = 1.06  $\times$  10<sup>16</sup>; (e and f)  $H_2/NO = 0.46$ ,  $NO_{In} = 0.90 \times 10^{16}$ .

according to  $(D1 - P7 - P5 + D4 - 1\frac{1}{2}P1)$ . The activation energies listed in Table 4 and Table 6. Rate parameters for  $NO-H<sub>2</sub>$  system satisfies the enthalpy balance.

In addition to the enthalpic constraint, there is also an entropic constraint. The former affects the activation energies while the latter affects the pre-exponential factors for simple Arrhenius type rate expressions. Application of an entropic balance requires that each of the constituent reaction step is reversible and follows the mass action kinetics. Thus, for lumped reaction steps that are irreversible or do not follow mass action kinetics, entropic constraints cannot be applied.

[Fig. 9](#page-10-0) shows the experimentally measured effect of the  $H_2/NO$ ratio on the product selectivity of  $N_2O$ ,  $N_2$ , and  $NH_3$  at 250 °C and 0.1 s delay time. The  $H_2/NO$  ratio of 2.26 results in the predominant formation of  $NH<sub>3</sub>$  while N<sub>2</sub> being a byproduct. A lower ratio of 1.04 results in the predominant  $N_2$  production and a small amount of  $NH<sub>3</sub>$  is also seen in the effluent. A further lower ratio of 0.46 results in formation of  $N_2$  and  $N_2O$ . The  $H_2$  is the limiting reactant so NO is also detected in the effluent.

The model comprising steps D1–D5 [\(Table 4](#page-7-0)), P1–P9 ([Table 6\)](#page-9-0) and A1–A2 [\(Table 8](#page-9-0)) predicts the predominant formation of  $NH<sub>3</sub>$ at  $H_2/NO$  ratio of 2.26 [\(Fig. 9](#page-10-0)a and b). The delayed appearance of  $H<sub>2</sub>O$  and NH<sub>3</sub> is well captured. The N<sub>2</sub> remains at a sustained level. Since NO is a limiting reactant, it is not seen in the effluent. The  $H_2$ / NO stoichiometry does not favor the formation of  $N_2O$ , and it is well explained by the simulations. The model predicts the delayed appearance of  $H_2$  in the effluent for excess of  $H_2$  case, but the experimental data show a sustained unreacted  $H_2$  in the exit. We do not have an explanation for this difference between model prediction and experiment.

The  $H<sub>2</sub>/NO$  ratio of 1.04 results in the predominant formation of N<sub>2</sub> as well as a very small amount of NH<sub>3</sub> also produced ([Fig. 9c](#page-10-0) and d). A smaller feed ratio of 0.46 leads to the formation of  $N_2$  and  $N_2O$ as well as unreacted NO ([Fig. 9e](#page-10-0) and f). The simulation predicts these trends. For example, the model predicts the stoichiometric formation of  $N_2$  for the initial pulses and a declining  $N_2$  production to a sustained value due to inhibition by surface oxygen. The model also predicts the breakthrough of NO and  $N_2O$ . The quantitative match for  $H_2/NO < 1$  is not as good because the  $NO/H_2$  rate parameters used in this study were taken from Xu et al. [\[9\]](#page-14-0) and were quantitatively good for  $H_2/NO > 1$ .

Temperature is an important parameter affecting the product selectivity of  $N_2$  and  $NH_3$ . Fig. 10 shows the experimentally observed effect of temperature for a fixed delay time of 0.1 s and H2/NO ratio of 2.26–2.55. Higher temperature facilitates N–O bond scission on reduced Pt surface, forming gaseous  $N<sub>2</sub>$  and leaving O on the surface. A subsequent  $H<sub>2</sub>$  pulse, after a delay time of 0.1 s, scavenges the surface O to form  $H_2O$ . In comparison, at lower temperature, the rate of N–O bond scission is lower than at 350  $\degree$ C. The higher interaction between NO or N with H leads to the increased formation of NH<sub>3</sub> at 250 °C.

[Fig. 11](#page-12-0) shows the experimentally measured effect of delay time on the selectivity of  $N_2$  and  $NH_3$  for a fixed temperature of 350 °C and  $H_2/NO$  ratio of 2.55. A shorter delay of ca. 0.0 s emulates a mixed feed in which NO and  $H_2$  simultaneously enter the catalyst zone. With such a feed protocol the interaction between surface NO, N, and H are more prominent, leading to formation of NH<sub>3</sub>. A longer delay between NO and  $H_2$ , e.g. 0.1 s, separates the NO and  $H<sub>2</sub>$  and the only way for these pulses to interact is through adsorbed O. As before, NO readily decomposes on Pt to form  $N_2$ ,



**Fig. 10.** Effect of temperature on product selectivity at 0.1 s delay time. (a and b)  $T = 250 \degree C$ , NO<sub>In</sub> = 1.09  $\times$  10<sup>16</sup>, H<sub>2</sub>/NO = 2.26; (c and d)  $T = 350 \degree C$ , NO<sub>In</sub> = 0.94  $\times$  10<sup>16</sup>,  $H_2/NO = 2.55$ .

<span id="page-12-0"></span>

**Fig. 11.** Effect of delay time on product selectivity at 350 °C. (a and b)  $\tau_d$  = 0.0 s, NO<sub>In</sub> = 0.94  $\times$  10<sup>16</sup>, H<sub>2</sub>/NO = 2.55; (c and d)  $\tau_d$  = 0.1 s, NO<sub>In</sub> = 0.94  $\times$  10<sup>16,</sup> H<sub>2</sub>/NO = 2.55.

leaving surface O on the Pt. The subsequent  $H_2$  pulse shows its presence on the Pt surface after a delay of 0.1 s. It scavenges the O from the surface, forming  $H_2O$ . As the delay time increases, the interaction between N and/or NO with H decreases. As a result,  $NH<sub>3</sub>$  formation decreases. The predominant route for  $N<sub>2</sub>$  formation for a longer delay time is solely due to NO decomposition. The formation of  $N_2$  and  $NH_3$  is a case of competition between N–N combination and combination of NO and/or N with H. The former produces  $N_2$  while the latter facilitates the formation of  $NH_3$ .

A closer look at the effect of delay time is demonstrated by plotting the exit fluxes and surface and gas concentration at the front of the catalyst zone for delay times 0.0 s, 0.1 s, and 1.0 s (Fig. 12).



Fig. 12. Exit fluxes, surface coverages and gaseous concentrations at front of catalyst zone for NO–H<sub>2</sub> pump–probe at 50th pulse, 350 °C, H<sub>2</sub>/NO = 2.5 and various delay times 0.0–1.0 s.



Fig. 13. Expanded view of Fig. 12. Exit fluxes, surface coverages and gaseous concentrations at front of catalyst zone for NO–H<sub>2</sub> pump–probe at 50th pulse, 350 °C, H<sub>2</sub>/ NO = 2.5 and various delay times 0.0–1.0 s.

An expanded view is provided in Fig. 13. The experimental conditions are 350 °C, H<sub>2</sub>/NO = 2.5, and NO<sub>In</sub> = 1.0  $\times$  10<sup>16</sup>. The exit flux, surface concentration, and gaseous concentration profiles are plotted for the 50th  $NO-H<sub>2</sub>$  pump–probe cycle. With a longer delay time,  $\tau_d$  = 1.0 s, the exit flux response shows formation of N<sub>2</sub> during the 0.0–1.0 s time interval. As seen in [Fig. 12,](#page-12-0) the corresponding Pt surface shows a sharp increase in the O coverage due to oxygen produced by NO decomposition. There is very small coverage of NO and N at the point of NO pulsing but their coverages decrease with time due to NO decomposition and  $N<sub>2</sub>$  formation, respectively. The O coverage is sustained for 1.0 s because of absence of surface oxygen removal in the form of  $O_2$ , NO<sub>2</sub>, etc. during that period.  $N<sub>2</sub>$  is the predominant gaseous species in the catalyst zone with traces of NO as depicted in the concentration profile plot. The catalyst zone shows a sustained level of  $H<sub>2</sub>O$  which slowly desorbs from the  $Al_2O_3$ .

A successive pulse of  $H_2$  at 1.0 s readily scavenges surface O as shown by the sudden decrease in oxygen coverage. The dominant species on the Pt is surface H. The concentration and flux profile also shows  $H_2$  as the dominant species after 1.0 s. One point to note is that the delay between incoming NO and  $H_2$  pulse is sufficiently long that NO and/or  $N_2$  species does not overlap with  $H_2$  in catalyst zone.

The exit flux and gas concentration profiles show that a smaller delay time of 0.1 s results in a larger overlap between N and H species. The  $N_2$  flux and concentration profiles for 0.1 s delay are sharper and narrower than those for the 1.0 s delay. This is due to an enhancement of  $N_2$  production due to surface reaction of surface NO and H (reaction step P4) in addition to NO decomposition. However, the integral amount of  $N_2$  formation for a single NO–H<sub>2</sub> pump probe cycle is almost the same for 0.1 s and 1.0 s delay.

A further reduction in delay time to 0.0 s shows an overlap of the  $H<sub>2</sub>$  pulse by the NO pulse. Since, the Knudsen diffusivity of  $H<sub>2</sub>$  is 3.9 times that of NO, it transports faster through the catalyst bed than NO. Although NO and  $H_2$  are fed simultaneously (more precisely,  $H_2$  is delayed by 1.2 ms),  $H_2$  overtakes NO in the inert zone and reaches the catalyst zone earlier than NO. The exit flux of  $N_2$  at 0.0 s delay is smaller than that of 0.1 s or 1.0 s delay. In addition to the presence of  $H<sub>2</sub>O$ , there is a sustained presence of  $NH<sub>3</sub>$  is seen in the catalyst zone. The surface coverage profile shows an interesting phenomenon: At the 0 s point, the H coverage shows a sharp peak due to  $H_2$  dissociation, followed by increase in the N and O coverages due to NO decomposition. After this, there is sustained H coverage due to  $NH<sub>3</sub>$  decomposition.

The surface coverage shows a longer residence time for H than N during the  $\tau_d$  to  $\tau_d$  +  $\tau_s$  interval. The product of N recombination is  $N_2$ , which is an inert species. Once  $N_2$  is formed, it does not readsorb on the Pt. On the other hand,  $H_2$  formed by H–H combination can dissociatively readsorb on the Pt. Hence, the coverages of H are higher than N during  $\tau_d$  to  $\tau_d + \tau_s$  period.

A comprehensive plot (Fig. 14) shows that  $NH<sub>3</sub>$  selectivity is a sensitive function of temperature and delay time between NO and  $H<sub>2</sub>$  pulses. A mixed feed with delay time 0.0 s shows the highest selectivity of  $NH<sub>3</sub>$  formation. As the delay time increases, while keeping the temperature constant, the formation of  $NH<sub>3</sub>$  decreases. A longer delay time translates to a less pronounced overlap between NO and  $H_2$  pulses on the catalyst, which leads to  $N_2$  formation via N–N combination. For a shorter delay, N and/or NO



Fig. 14. NH<sub>3</sub> selectivity as function of temperature and delay times for H<sub>2</sub>/NO = 2.0– 2.5.

<span id="page-14-0"></span>combines with surface H to form NH<sub>3</sub>. Higher temperature also facilitates  $N<sub>2</sub>$  formation since, NO decomposition rates are higher and so does the N–N combination.

#### 5. Conclusions

We have presented a comprehensive multi-pulse microkinetic model for NO decomposition and NO–H<sub>2</sub> pump–probe over Pt/  $Al_2O_3$  catalyst that explains the essential experimental observations. NO pulsing experiments show the inhibition effect of oxygen poisoning in the studied temperature range. The  $NO-H<sub>2</sub>$  pump– probe experiments demonstrate the effect of temperature,  $H_2/NO$ ratio ( $\geq$ 1), and delay time between NO and H<sub>2</sub> pulses on N<sub>2</sub>,  $N_2$ O, and  $NH_3$  selectivity. The higher temperature favors NO decomposition to form  $N_2$ . The higher ratio of  $H_2/NO$  favors  $NH_3$ production. A longer delay time between NO and  $H_2$  would form  $N_2$  as the dominant product while shorter delay time facilitates interaction between N and/or NO with H to form  $NH<sub>3</sub>$ . The model also accounts for the physical adsorption/desorption on  $Al_2O_3$ . The model captures the competition between surface N–N recombination and N and/or NO reaction with H to form  $N_2$  and NH<sub>3</sub>, respectively.

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